

A semi-Lagrangian AMR scheme for 2D transport problems in conservation form

Pep Mulet Mestre, Francesco Vecil

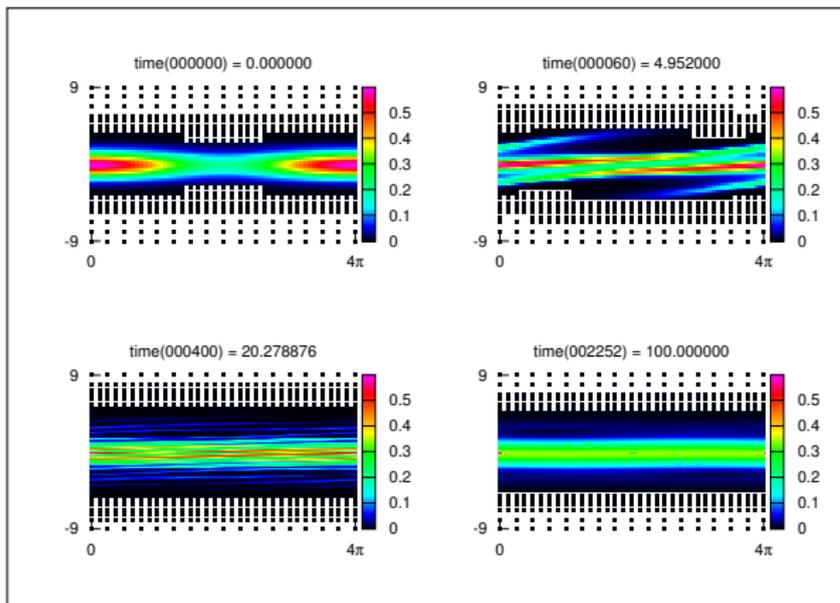
Universitat de València

CSASC, Koper/Capodistria (Slovenia), 12 June 2013

Outline

- 1 Introduction
- 2 Numerical tools
 - Multiresolution framework
 - Time integration
- 3 Experiments
 - Introduction
 - 1D tests
 - 2D tests

Motivation



No need for fine meshing everywhere in the domain.



Refine only where the important information is.

Framework

Equations

In dimension N , transport equations written in conservation form:

$$\frac{\partial u}{\partial t} + \frac{\partial}{\partial x_1} [a_1(t, x_1, x_2)u] + \frac{\partial}{\partial x_2} [a_2(t, x_1, x_2)u] = 0, \quad u(0, x_1, x_2) = u^0(x_1, x_2),$$

where $\mathbf{a} : \mathbb{R}_{\geq 0} \times \Omega \rightarrow \mathbb{R}^2$ is the advection field.

Example

The three-dimensional Vlasov-Maxwell equation

$$\frac{\partial f}{\partial t} + \mathbf{v}(\mathbf{p}) \cdot \frac{\partial f}{\partial \mathbf{x}} + \mathbf{F} \cdot \frac{\partial f}{\partial \mathbf{p}} = 0, \quad \mathbf{v}(\mathbf{p}) := \frac{\mathbf{p}}{m\sqrt{1 + \frac{|\mathbf{p}|^2}{m^2 c^2}}}, \quad \mathbf{F} := -e(\mathbf{E} + \mathbf{v}(\mathbf{p}) \wedge \mathbf{B}),$$

describes the evolution of $f(t, \mathbf{x}, \mathbf{p})$, typically representing the concentration of electrons or holes at position \mathbf{x} and momentum \mathbf{p} .

Features

Shocks, large gradients, filamentation, microscopic structures.

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Grid hierarchy

Resolution levels

We define $L + 1$ resolution levels: the coarsest is $\ell = 0$, the finest $\ell = L$. In 1D, the meshes are

$$x_{\ell,j} = x_{\min} + j\Delta x_{\ell}, \quad \Delta x_{\ell} = \frac{x_{\max} - x_{\min}}{2^{\ell}N_0}.$$

Grid

The ℓ -grid at time t^n is

$$G_{\ell}^n = \{x_{\ell,j}\}_{j \in \mathcal{G}_{\ell}^n}.$$

We are interested in

$$\mathcal{G}_{\ell}^n \subseteq \prod_{i=1}^N \{0, \dots, N_{i,\ell}\}.$$

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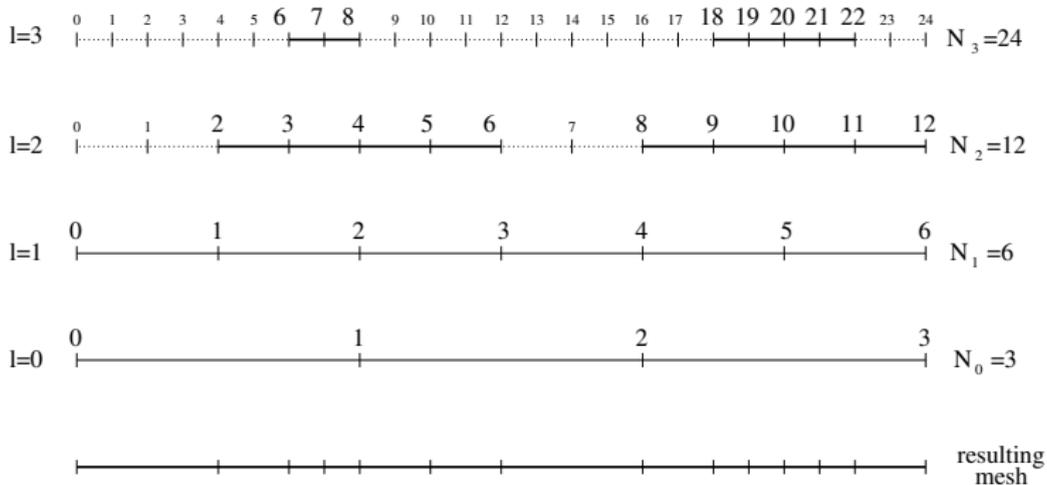
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Nesting condition

We are interested in *nested* meshes:



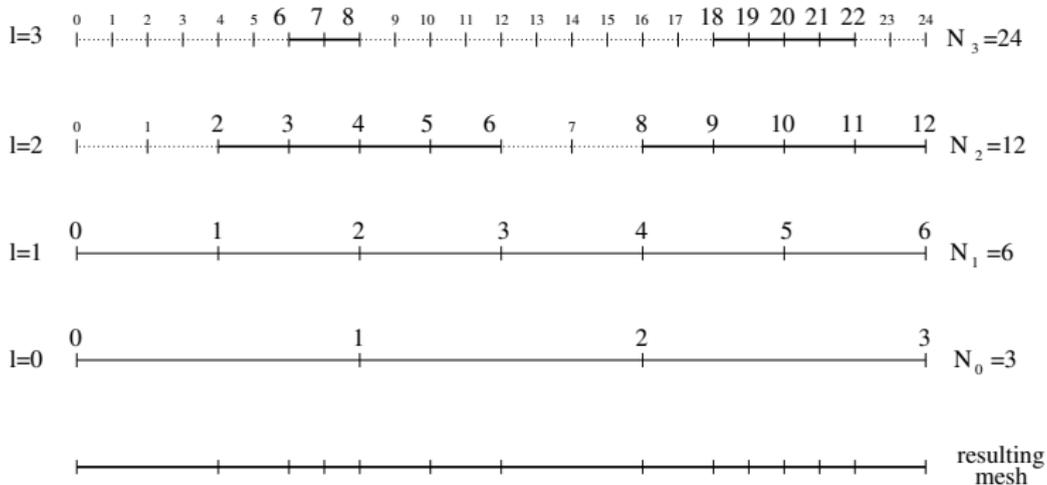
Ghost points

Ghost points are added outside the l -grids to take into account the boundary conditions for the time integration.

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Grid selection

The grid selection is based on two criteria: *prediction* and *gradient*.

The *prediction* criterion

Let $\{u_{\ell,j}\}_{j \in \mathcal{G}_\ell^n}$ the point values at resolution level ℓ . Let $p[u_\ell]$ an interpolator at resolution level ℓ . We keep point $x_{\ell,j}$ for refinement if

$$|p[u_{\ell-1}](x_{\ell,j}) - u_{\ell,j}| > \tau_p.$$

The *gradient* criterion

We estimate the gradient at point $x_{\ell,j}$ by means of the discrete gradient computed at resolution level $\ell - 1$: if it is over a certain tolerance parameter $\tau_{d,\ell-1}$, then point $x_{\ell,j}$ is selected for refinement.

Reconstruction

Once the grid has been selected, reconstruct by means of an interpolator the point values that are not assigned yet.

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1D semi-Lagrangian strategy

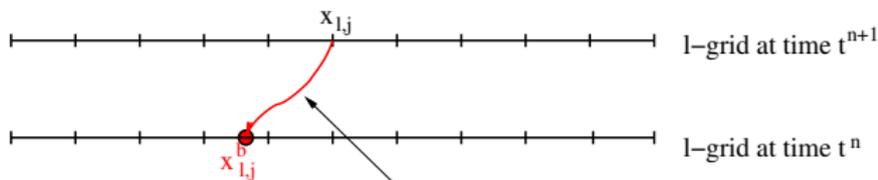
Characteristic-based solution

The solution to the PDE
$$\frac{\partial u}{\partial t} + \frac{\partial}{\partial x} [a(t, x)u] = 0, \quad u(0, x) = u^0(x)$$

is given by
$$u(t, x) = u(s, \mathcal{X}(s; t, x)) J(s; t, x),$$

with $\mathcal{X}(s; t, x)$ the characteristic at time s , starting from x at time t :

$$\frac{d\mathcal{X}(s; t, x)}{ds} = a(s, \mathcal{X}(s; t, x)), \quad \mathcal{X}(t; t, x) = x, \quad J(s; t, x) := \frac{\partial \mathcal{X}(s; t, x)}{\partial x}.$$



the characteristics
are solved backward
from time t^{n+1} to time t^n
through the advection field $a(t, x)$

1D semi-Lagrangian strategy

What do we need?

- A solver for the characteristics $\mathcal{X}(s; t, x)$: **Runge-Kutta**.
- An approximation for the Jacobian through **centred finite differences**:

$$J(s; t, x) = \frac{\partial \mathcal{X}(s; t, x)}{\partial x} \approx \frac{\tilde{\mathcal{X}}(s; t, x + \delta x) - \tilde{\mathcal{X}}(s; t, x - \delta x)}{2\delta x}, \quad \delta x = 10^{-m} \Delta x_\ell.$$

- An interpolator to reconstruct $u(s, \mathcal{X}(s; t, x))$: **PVWENO** (Point-Value Weighted Essentially Non-Oscillatory).

Error estimate

The local truncation error can be estimated

$$E = \underbrace{\mathcal{O}(\Delta t^{s+1})}_{\text{Runge-Kutta}} + \underbrace{\mathcal{O}(\Delta x_\ell^2)}_{\text{Jacobian}} + \underbrace{\mathcal{O}(\Delta x_\ell^{2r})}_{\text{PVWENO (advection)}}.$$

approximation of the characteristics

1D semi-Lagrangian strategy

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The 2D case

Grid hierarchy and selection

We do not give details, but we apply strategies similar to the 1D case.

The 2D PDE

We solve the 2D PDE
$$\frac{\partial u}{\partial t} + \frac{\partial}{\partial x_1} (a_1 u) + \frac{\partial}{\partial x_2} (a_2 u) = 0$$

by splitting the (x_1, x_2) -domain thanks to the second-order Strang scheme:

- Solve for a $\frac{\Delta t}{2}$ time step $\frac{\partial u}{\partial t} + \frac{\partial}{\partial x_1} (a_1 u) = 0$;
- Solve for a Δt time step $\frac{\partial u}{\partial t} + \frac{\partial}{\partial x_2} (a_2 u) = 0$;
- Solve for a $\frac{\Delta t}{2}$ time step $\frac{\partial u}{\partial t} + \frac{\partial}{\partial x_1} (a_1 u) = 0$.

Error estimate

The Strang splitting constrains the accuracy:

$$E = \mathcal{O}\left(\Delta t^{\min(s+1,3)}\right) + \underbrace{\mathcal{O}(\Delta x_{1,\ell}^2) + \mathcal{O}(\Delta x_{2,\ell}^2)}_{\text{approximated Jacobian}} + \underbrace{\mathcal{O}(\Delta x_{1,\ell}^{2r}) + \mathcal{O}(\Delta x_{2,\ell}^{2r})}_{\text{PVWENO}}.$$

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Goals

The Adaptive-Mesh-Refinement (AMR) framework is compared to the equivalent Fixed-Mesh (FM) results.

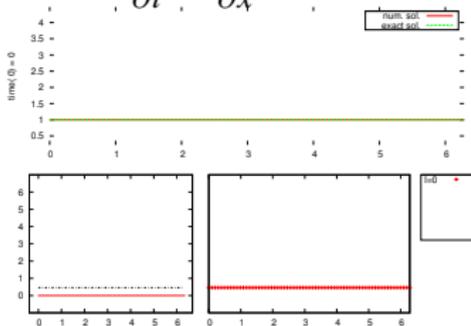
Of course, AMR cannot be more accurate than FM. Rather, it achieves faster computational times in exchange of a loss of precision.

Outline

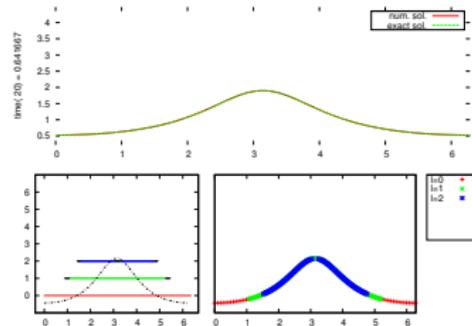
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Variable-coefficient advection

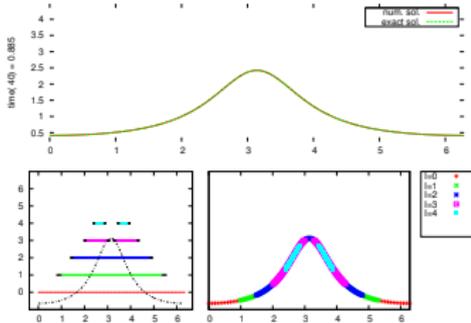
The test case $\frac{\partial u}{\partial t} + \frac{\partial}{\partial x} (\sin(x) u) = 0$ produces a blow-up.



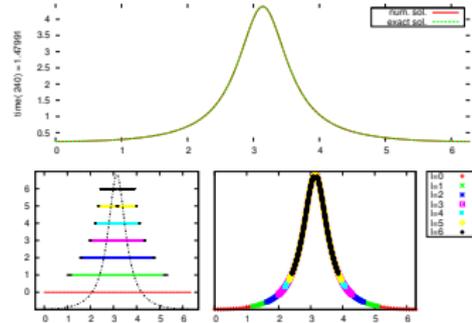
(a) at time 0



(b) at time 0.65



(c) at time 0.89



(d) at time 1.48

Variable-coefficient advection

Speedup

For parameters

$N_0 = 128$	points at $\ell = 0$
$L = 4$	number of resolution levels
$\Delta t_0 = 0.125$	maximum time step
$\tau_p = 10^{-4}$	<i>prediction</i> -criterion tolerance
$\tau_{d,0} = 0.5$	<i>gradient</i> -criterion tolerance,

AMR reaches a speedup of 35 times with respect to the equivalently-resolved FM, with a loss of precision from 10^{-9} to roughly 10^{-6} (the L^2 -error w.r.t. the analytical solution).

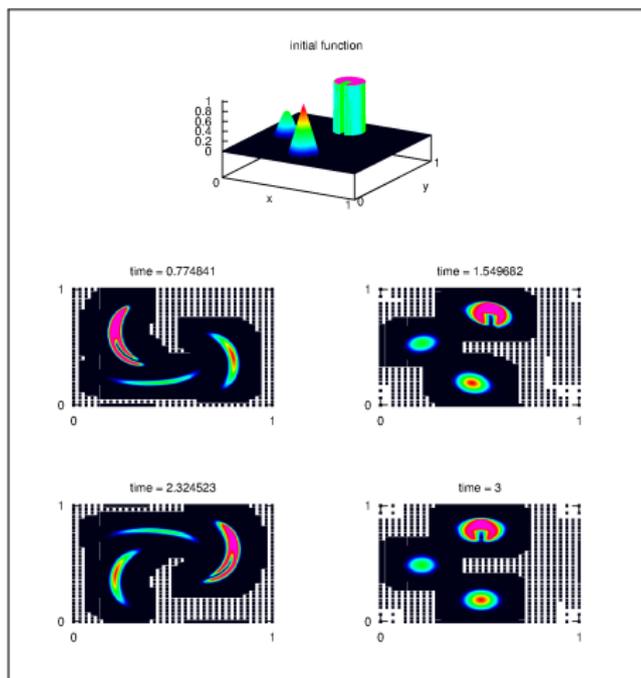
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Deformation flows

The system

$$\frac{\partial f}{\partial t} + \frac{\partial}{\partial x} \left[\sin^2(\pi x) \sin(2\pi y) g(t) f \right] + \frac{\partial}{\partial y} \left[-\sin^2(\pi y) \sin(2\pi x) g(t) f \right] = 0, \quad (x, y) \in [0, 1]^2,$$



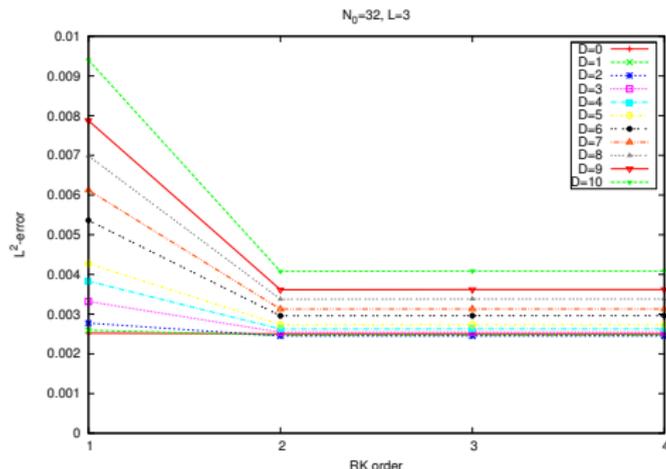
for $g(t) = \cos\left(\frac{\pi t}{T}\right)$,
 periodically recovers the
 initial datum after
 alternate clockwise and
 counterclockwise
 twistings.

Performances

Speedup ≈ 2 .

Deformation flows

The ODE integrator for the characteristics



As announced by the error estimate

$$E = \mathcal{O}(\Delta x_{1,\ell}^2) + \mathcal{O}(\Delta x_{2,\ell}^2) + \mathcal{O}\left(\Delta t^{\min(s+1,3)}\right)$$

the Strang-splitting order constrains the accuracy.

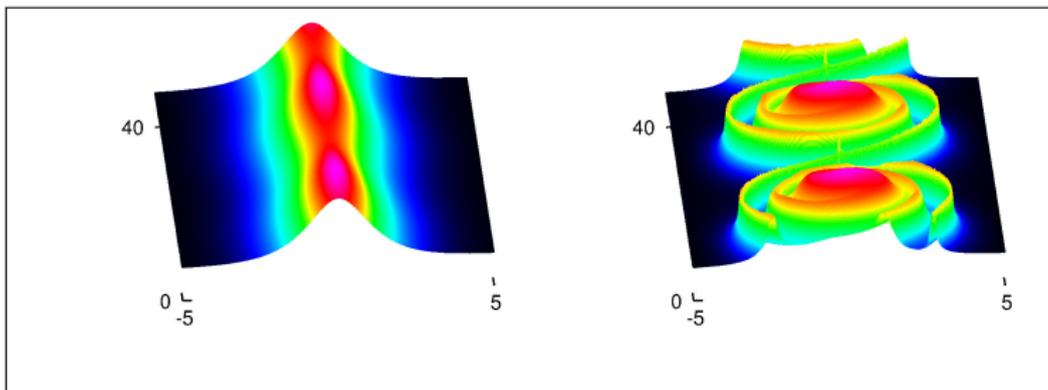
Kelvin-Helmholtz instabilities

The model

The *guiding-center* model (omitting some details)

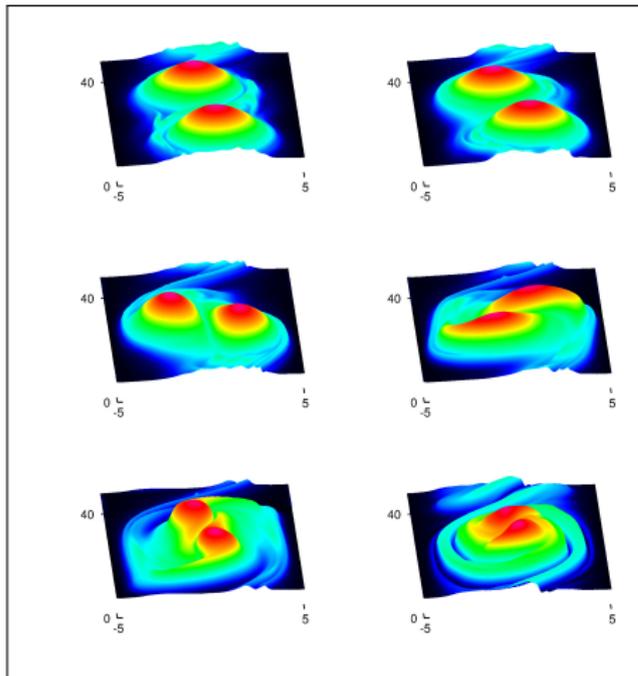
$$\frac{\partial \rho}{\partial t} + \frac{\partial}{\partial x_1} \left[\frac{\partial \Phi}{\partial x_2} \rho \right] + \frac{\partial}{\partial x_2} \left[-\frac{\partial \Phi}{\partial x_1} \rho \right] = 0, \quad \Delta_{x_1, x_2} \Phi = \rho,$$

for initial condition $\rho(0, x_1, x_2) = 1.5 \operatorname{sech} \left(\frac{x_2}{0.9} \right) \cdot (1 + 0.08 \sin(2kx_1))$,
 periodic x_1 - and Dirichlet x_2 -boundaries, produces vortices and filamentation.

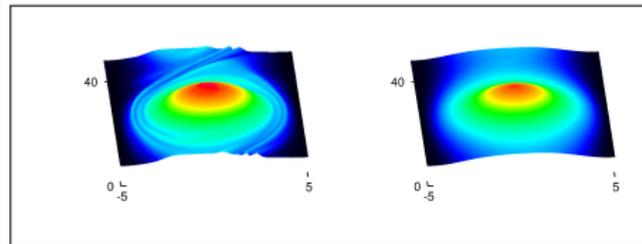


(f) first instability

Kelvin-Helmholtz instabilities



(g) second instability



(h) steady state

SPEEDUP

The AMR strategy achieves a speedup of about 2.8, for parameters $(N_{*,0}, L) = (32, 4)$ (the maximum resolution is 512×512), $\tau_p = 10^{-3}$, $\tau_{d,0} = 1.5$.

GRÀCIES!
HVALA!
DANKE!
ĎAKUJEM!
DĚKUJI!